

Solving PDEs with PGI CUDA Fortran

Part 6: More methods for more partial differential equations

Outline

Heat equation in 1D: implicit and Crank-Nicolson schemes. Heat equation in more dimensions: alternating-direction implicit method. Multigrid method. Wave equation in 1D and 2D: strings and drums.

Heat equation in 1D: more schemes

A symbol for the difference operator

$$\delta_x^2 u_j^n \equiv u_{j-1}^n - 2u_j^n + u_{j+1}^n$$

FTCS scheme with Dirichlet boundary conditions

$$u_j^{n+1} = u_j^n + \beta \delta_x^2 u_j^n, \quad \beta = dt/dx^2$$

$$u_j^{n+1} = u_j^n + \beta (u_{j-1}^n - 2\beta u_j^n + u_{j+1}^n)$$

Features: 1st-order accurate in time, 2nd-order in space, conditionally stable ($\beta \leq 1/2$)

BTCS scheme (backward-time centered-space)

implicit formula

$$u_j^{n+1} = u_j^n + \beta \delta_x^2 u_j^{n+1}, \quad \beta = dt/dx^2$$

$$\begin{pmatrix} 1+2\beta & -\beta & & & \\ -\beta & 1+2\beta & -\beta & & \\ & & \cdot & & \\ & & & -\beta & 1+2\beta & -\beta \\ & & & -\beta & 1+2\beta \end{pmatrix} \begin{pmatrix} u_1^{n+1} \\ u_2^{n+1} \\ \cdot \\ u_{j-1}^{n+1} \\ u_j^{n+1} \end{pmatrix} = \begin{pmatrix} u_1^n + \beta u_0 \\ u_2^n \\ \cdot \\ u_{j-1}^n \\ u_j^n + \beta u_{j+1} \end{pmatrix}$$

Features: 1st-order accurate in time, 2nd-order in space, unconditionally stable (i.e., for any dt)

Each time step requires direct solution to a linear algebraic system with tridiagonal matrix of size $J \times J$.

Crank-Nicolson scheme (CN)

implicit formula with an average of FTCS and BTCS schemes on the right-hand side

$$u_j^{n+1} = u_j^n + \frac{\beta}{2} (\delta_x^2 u_j^n + \delta_x^2 u_j^{n+1})$$

$$\begin{pmatrix} 1+\beta & -\beta/2 & & & \\ -\beta/2 & 1+\beta & -\beta/2 & & \\ & & \cdot & & \\ & & & -\beta/2 & 1+\beta \end{pmatrix} \begin{pmatrix} u_1^{n+1} \\ u_2^{n+1} \\ \cdot \\ u_{j-1}^{n+1} \\ u_j^{n+1} \end{pmatrix} = \begin{pmatrix} 1-\beta & \beta/2 & & & \\ \beta/2 & 1-\beta & \beta/2 & & \\ & & \cdot & & \\ & & & \beta/2 & 1-\beta \end{pmatrix} \begin{pmatrix} u_1^n \\ u_2^n \\ \cdot \\ u_{j-1}^n \\ u_j^n \end{pmatrix} + \begin{pmatrix} \beta u_0 \\ 0 \\ \cdot \\ \beta u_{j+1} \end{pmatrix}$$

Features: 2nd-order accurate in both time and space, unconditionally stable

Each time step requires direct solution to a linear algebraic system with tridiagonal matrix of size $J \times J$.

Heat equation in 2D: FTCS, BTCS and CN schemes

Difference operators $\delta_x^2 u_{jk}^n \equiv u_{j-1,k}^n - 2u_{jk}^n + u_{j+1,k}^n$, $\delta_y^2 u_{jk}^n \equiv u_{j,k-1}^n - 2u_{jk}^n + u_{j,k+1}^n$

FTCS scheme $u_{jk}^{n+1} = u_{jk}^n + \beta (\delta_x^2 u_{jk}^n + \delta_y^2 u_{jk}^n)$

BTCS scheme $u_{jk}^{n+1} = u_{jk}^n + \beta (\delta_x^2 u_{jk}^{n+1} + \delta_y^2 u_{jk}^{n+1})$

CN scheme $u_{jk}^{n+1} = u_{jk}^n + \frac{\beta}{2} (\delta_x^2 u_{jk}^n + \delta_y^2 u_{jk}^n + \delta_x^2 u_{jk}^{n+1} + \delta_y^2 u_{jk}^{n+1})$

For implicit BTCS and CN schemes, the matrix is $J^2 \times J^2$, sparse and band diagonal (tridiagonal with fringes).

Direct solution is possible with special methods.

Heat equation in more dimensions: alternating-direction implicit (ADI) method

2D: splitting the time step into 2 substeps, each of length $t/2$

$$\begin{aligned} u_{jk}^{n+1/2} &= u_{jk}^n + \frac{\beta}{2} \left(\delta_x^2 u_{jk}^{n+1/2} + \delta_y^2 u_{jk}^n \right) \\ u_{jk}^{n+1} &= u_{jk}^{n+1/2} + \frac{\beta}{2} \left(\delta_x^2 u_{jk}^{n+1/2} + \delta_y^2 u_{jk}^{n+1} \right) \end{aligned}$$

3D: splitting the time step into 3 substeps, each of length $t/3$

$$\begin{aligned} u_{jkl}^{n+1/3} &= u_{jkl}^n + \frac{\beta}{3} \left(\delta_x^2 u_{jkl}^{n+1/3} + \delta_y^2 u_{jkl}^n + \delta_z^2 u_{jkl}^n \right) \\ u_{jkl}^{n+2/3} &= u_{jkl}^{n+1/3} + \frac{\beta}{3} \left(\delta_x^2 u_{jkl}^{n+1/3} + \delta_y^2 u_{jkl}^{n+2/3} + \delta_z^2 u_{jkl}^{n+1/3} \right) \\ u_{jkl}^{n+1} &= u_{jkl}^{n+2/3} + \frac{\beta}{3} \left(\delta_x^2 u_{jkl}^{n+2/3} + \delta_y^2 u_{jkl}^{n+2/3} + \delta_z^2 u_{jkl}^{n+1} \right) \end{aligned}$$

All substeps are implicit and each requires direct solutions to J independent linear algebraic systems with tridiagonal matrices of size $J \times J$.

Example: ADI method for heat equation in 2D and 3D

Wave equation

a quantity travelling over the domain

a partial differential equation (2nd-order in time t , 2nd-order in spatial variables X) for a function $u(t, X)$

1D (one-dimensional) case: $X = x$, 2D case: $X = x, y$, 3D case: $X = x, y, z$

General form: $\partial_t^2 u(t, X) = c^2 \Delta u(t, X)$
 in 3D: $\partial_t^2 u(t, x, y, z) = c^2 (\partial_x^2 + \partial_y^2 + \partial_z^2) u(t, x, y, z)$
 Initial conditions: $u(t_0, X) = u_0(X), \quad \partial_t u(t_0, X) = v_0(X)$
 Boundary conditions: $u(t, X_B) = u_B(t, X_B)$ on the boundary
 i.e., the **initial value problem (IVP)** for the **hyperbolic partial differential equation**

Discretized wave equation in 1D

1D wave equation

$$\partial_t^2 u(t, x) = c^2 \partial_x^2 u(t, x)$$

can be rewritten into the form of two equations of the 1st-order in time

$$\begin{aligned} (\partial_t + c \partial_x) u(t, x) &= v(t, x) \\ (\partial_t - c \partial_x) v(t, x) &= 0 \end{aligned}$$

Discretization grids

$$\begin{aligned} t_n &= t_0 + n dt, & dt &= (t_N - t_0)/N \\ x_j &= x_0 + j dx, & dx &= (x_J - x_0)/J \\ u_j^n &\approx u(t_n, x_j), & v_j^n &\approx v(t_n, x_j) \end{aligned}$$

Explicit FTBS scheme (forward-in-time, backward-in-space)

FD1 for time: $\partial_t u_j^n \approx (u_j^{n+1} - u_j^n)/dt$
 FD1 for space: $\partial_x u_j^n \approx (-u_{j-1}^n + u_j^n)/dx$

Features: low accuracy, **stability** for $\frac{c dt}{dx} \leq 1$ (**Courant-Friedrichs-Lewy condition**)

PDEs in the matrix form:

$$\partial_t \begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} -c \partial_x & 1 \\ 0 & c \partial_x \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

Discretized equations:

$$\begin{pmatrix} u \\ v \end{pmatrix}^{n+1} = \begin{pmatrix} I - \gamma A & \delta I \\ 0 & I + \gamma A \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}^n, \quad \gamma = \frac{c dt}{dx}, \quad \delta = dt, \quad A = \begin{pmatrix} 1 & 0 & 0 & \cdot \\ -1 & 1 & 0 & \cdot \\ 0 & -1 & 1 & \cdot \end{pmatrix}$$

and I is the identical matrix

Explicit FTCS scheme (forward-in-time, centered-in-space)

FD1 for time: $\partial_t u_j^n \approx (u_j^{n+1} - u_j^n)/dt$

FD2 for space: $\partial_x u_j^n \approx (-u_{j-1}^n + u_{j+1}^n)/(2dx)$

Features: **unstable** for any dt, i.e., FTCS scheme inappropriate for the wave equation

Implicit Crank-Nicolson scheme

implicit formula with an average of FTBS and BTBS schemes on the right-hand side

$$\begin{pmatrix} u \\ v \end{pmatrix}^{n+1} = \begin{pmatrix} u \\ v \end{pmatrix}^n + \begin{pmatrix} -\gamma A & \delta I \\ 0 & \gamma A \end{pmatrix} \left[\begin{pmatrix} u \\ v \end{pmatrix}^n + \begin{pmatrix} u \\ v \end{pmatrix}^{n+1} \right], \quad \gamma = \frac{c dt}{dx}, \quad \delta = dt, \quad A = \begin{pmatrix} 1 & 0 & 0 & \cdot \\ -1 & 1 & 0 & \cdot \\ 0 & -1 & 1 & \cdot \end{pmatrix}$$

Features: higher accuracy, **unconditional stability** (i.e., for any dt)

Example: travelling waves

domain $t \geq t_0 = 0, x \geq x_0 = 0$

initial condition $u(0, x) = u_0(x), v(0, x) = 0$

boundary condition $u(t, 0) = 0$

analytical solution $u(t, x) = u_0(t - cx)$

Links and references

PDEs

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Chapter 19.2: Diffusive initial value problems

Chapter 19.3: Initial value problems in multidimensions

Chapter 19.6: Multigrid methods for boundary value problems

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Wave equation on GPU

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